

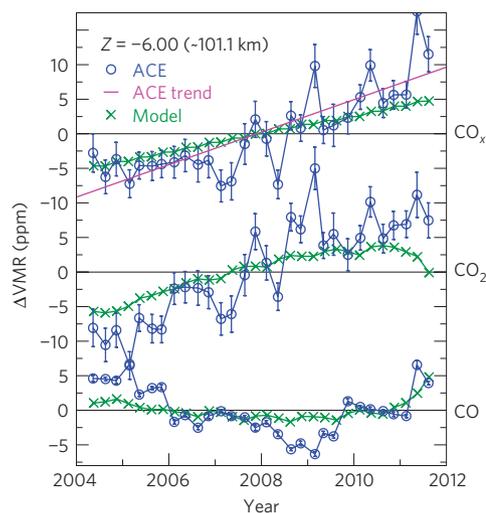
# Observations of increasing carbon dioxide concentration in Earth's thermosphere

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Carbon dioxide occurs naturally throughout Earth's atmosphere. In the thermosphere, CO<sub>2</sub> is the primary radiative cooling agent and fundamentally affects the energy balance and temperature of this high-altitude atmospheric layer<sup>1,2</sup>. Anthropogenic CO<sub>2</sub> increases are expected to propagate upward throughout the entire atmosphere, which should result in a cooler, more contracted thermosphere<sup>3–5</sup>. This contraction, in turn, will reduce atmospheric drag on satellites and may have adverse consequences for the orbital debris environment that is already unstable<sup>6,7</sup>. However, observed thermospheric mass density trends derived from satellite orbits are generally stronger than model predictions<sup>8,9</sup>, indicating that our quantitative understanding of these changes is incomplete. So far, CO<sub>2</sub> trends have been measured only up to 35 km altitude<sup>10–12</sup>. Here, we present direct evidence that CO<sub>2</sub> concentrations in the upper atmosphere—probably the primary driver of long-term thermospheric trends—are increasing. We analyse eight years of CO<sub>2</sub> and carbon monoxide mixing ratios derived from satellite-based solar occultation spectra. After correcting for seasonal–latitudinal and solar influences, we obtain an estimated global increase in CO<sub>x</sub> (CO<sub>2</sub> and CO, combined) concentrations of  $23.5 \pm 6.3$  ppm per decade at an altitude of 101 km, about 10 ppm per decade faster than predicted by an upper atmospheric model. We suggest that this discrepancy may explain why the thermospheric density decrease is stronger than expected.

In the mesosphere (50–90 km altitude) and thermosphere (>90 km), almost all carbon is partitioned between CO<sub>2</sub> and CO through ultraviolet photolysis of CO<sub>2</sub> to create CO and chemical loss of CO to re-form CO<sub>2</sub> (refs 1,13). The photochemistry and atmospheric dynamics introduce solar cycle, seasonal and latitudinal dependences into the partitioning and the overall carbon abundance<sup>13,14</sup>. The CO<sub>2</sub> volume mixing ratio (VMR) is approximately constant up to 65–80 km and falls off rapidly at higher altitudes owing to molecular diffusion and photolysis<sup>13</sup>. The proportion of carbon in the form of CO increases from less than 3% below 80 km to more than 20% above 100 km (refs 13,14). CO<sub>2</sub> cools the mesosphere and thermosphere through collisional excitation (by atomic oxygen) of its bending vibrational mode and subsequent 15 μm band emission, most of which escapes from the thin upper atmosphere<sup>12</sup>.

We analysed temporal changes in CO<sub>2</sub> and CO VMRs measured by the Atmospheric Chemistry Experiment Fourier Transform Spectrometer<sup>15</sup> (ACE-FTS) from April 2004 to September 2011, using the approach described in Methods. Figure 1 shows the CO<sub>2</sub>



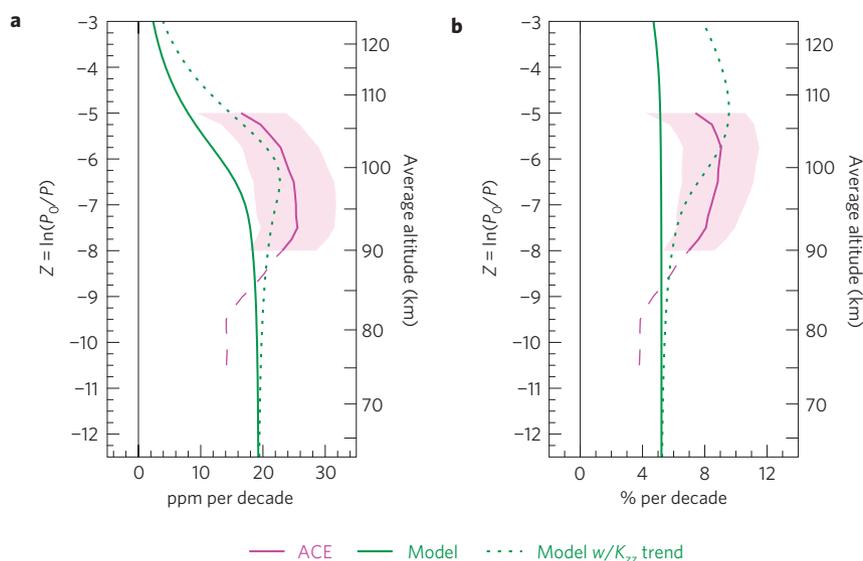
**Figure 1** | Residual VMRs of CO, CO<sub>2</sub> and CO<sub>x</sub> = CO + CO<sub>2</sub>, after removal of seasonal–latitudinal effects. Results are shown at pressure level  $Z = -6$  (altitude  $\sim 101$  km) for ACE-FTS (blue circles) and the NCAR global mean model (green crosses). The solid pink line is the linear trend from a least-squares fit of the ACE-FTS CO<sub>x</sub> residuals. Error bars denote the estimated  $1\sigma$  uncertainty of the mean.

and CO residual time series, after removal of seasonal–latitudinal variations and the overall mean, at a pressure level corresponding to an altitude of  $\sim 101$  km. Also shown is a corresponding analysis of the combined CO<sub>2</sub> and CO VMRs, which we term CO<sub>x</sub>. As expected from photochemistry, the CO VMR varies in phase with the solar cycle (with a minimum in 2008–2009), whereas the CO<sub>2</sub> increase is attenuated during the 2009–2012 ascending phase of the solar cycle. The CO<sub>x</sub> VMR, however, increases fairly steadily from 2004 to 2012, with a least-squares linear trend of  $23.5 \pm 6.3$  ppm per decade, where the uncertainty is the 95% confidence interval. As the partitioning of carbon between CO and CO<sub>2</sub> is influenced primarily by solar ultraviolet irradiance, the trend in CO<sub>x</sub> should be largely free of solar cycle variations and thus can be used to estimate the long-term relative trend in CO<sub>2</sub>, assuming the relative proportion of CO and CO<sub>2</sub> is stable.

For comparison with the ACE-FTS trends, we conducted simulations with the National Center for Atmospheric Research (NCAR) global mean model of the mesosphere, thermosphere and ionosphere<sup>1</sup>. We conducted a 1992–2012 global mean model

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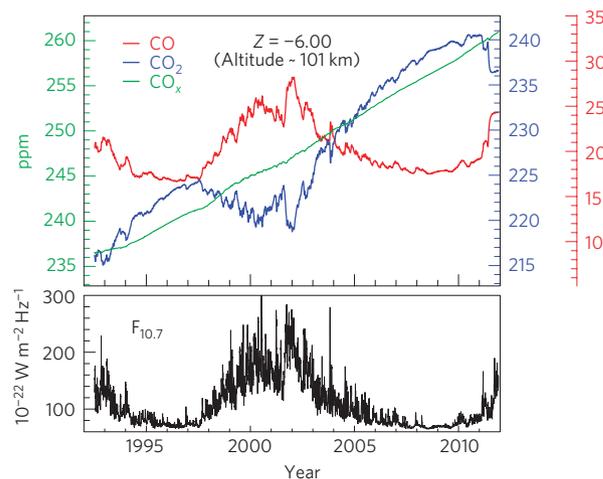
**Figure 2 | Height dependence of 2004–2012  $\text{CO}_x$  linear trends.** **a**, Absolute trends. **b**, Relative trends, obtained by dividing the absolute trends by the profiles shown in Fig. 4. Results are shown for ACE-FTS (pink), the NCAR global mean model (solid green) and the global mean model with a 15% per decade eddy diffusion trend added (dotted green). The shaded areas encompass the estimated 95% confidence interval of the ACE-FTS trends. The dashed lines denote the portion of the ACE-FTS trend profiles that is influenced by the prescribed stratospheric  $\text{CO}_2$  trends.

simulation using interpolated annual global mean  $\text{CO}_2$  VMRs measured in the troposphere<sup>16</sup> to specify  $\text{CO}_2$  variations at the model's lower boundary of 30 km, and we repeated the ACE-FTS analysis procedure on the 2004.25–2012.75 daily output of the model simulation. The model  $\text{CO}_x$  residual time series is shown in Fig. 1, and exhibits a noticeably weaker trend ( $13.2 \pm 0.3$  ppm per decade) than the ACE-FTS data. Figure 2a shows height profiles (computed on a pressure grid as described in Methods) of the  $\text{CO}_x$  linear trends. The ACE-FTS trends between 90 and 105 km decrease with increasing altitude at approximately the same rate as the model trends, but the ACE-FTS trends are 5–10 ppm per decade larger than the model trends. Figure 2b shows the relative trends (that is, normalized to the average  $\text{CO}_x$  VMR), which are approximately constant with height between 90 and 105 km:  $\sim 8\%$  per decade for ACE-FTS,  $\sim 5\%$  per decade for the global mean model.

In Fig. 2, the decrease in the ACE-FTS trends below 90 km is probably an artefact of the retrieval process, which specifies the  $\text{CO}_2$  VMR a priori below  $\sim 70$  km and transitions to free retrieval at higher altitudes<sup>17</sup>. The imposed stratospheric trend of 15 ppm per decade is based on 1980s and 1990s tropospheric  $\text{CO}_2$  data; the observed tropospheric trend averaged  $\sim 20$  ppm per decade during the 2000s. We tested the effect of using the more realistic stratospheric trend by reprocessing a subset of the ACE-FTS occultations; the measured trend above 90 km changed by less than 1.5 ppm per decade, which is well within the statistical uncertainty of 6.3 ppm per decade. We also fitted the global mean model  $\text{CO}_2$  profiles to the ACE-FTS  $\text{CO}_2$  retrieval function and found that the model trends above 90 km are uninfluenced by the stratospheric constraint.

Figure 3 shows the 1992–2012 temporal variation of  $\text{CO}_2$ , CO and  $\text{CO}_x$  VMRs from the model simulation at 101 km. The solar cycle amplitude of  $\text{CO}_x$  is less than 1 ppm, which is much smaller than the expected decadal  $\text{CO}_2$  increase of  $> 10$  ppm and supports our use of  $\text{CO}_x$  trends as a proxy for long-term  $\text{CO}_2$  trends. However, Fig. 1 indicates that the solar cycle dependence is stronger in the ACE-FTS data than in the model, so the model's representation of solar influences is probably incomplete.

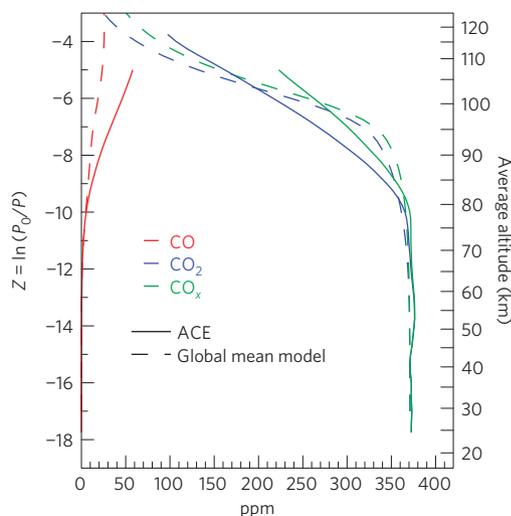
Figure 4 shows average absolute profiles of  $\text{CO}_2$ , CO and  $\text{CO}_x$  VMRs from the model and from ACE-FTS for 2004.25–2011.75.



**Figure 3 | Temporal variation of carbon at pressure level  $Z = -6$  (altitude  $\sim 101$  km) from the NCAR global mean model simulation.** Shown are VMRs of CO (red),  $\text{CO}_2$  (blue) and  $\text{CO}_x = \text{CO} + \text{CO}_2$  (green). The data are plotted according to the colour-coordinated y axes. The bottom panel shows the 10.7 cm solar radio flux ( $F_{10.7}$ ), a proxy for solar ultraviolet irradiance.

The  $\text{CO}_2$  VMR begins to fall off above  $\sim 80$  km in the data and above  $\sim 95$  km in the model. There are long-standing discrepancies between measured and modelled  $\text{CO}_2$  profiles, particularly with respect to the height of departure from a well-mixed state (the knee; refs 13,14). Models also generally underestimate the CO VMR in the mesosphere and lower thermosphere<sup>13</sup>, although good agreement was obtained<sup>14</sup> between the ACE-FTS data and the Canadian Middle Atmosphere Model. Our comparisons (Figs 1 and 2) between data and model temporal changes should not be significantly affected by the discrepancies in the baseline profiles.

The 10 ppm per decade difference between the ACE-FTS and global mean model  $\text{CO}_x$  trends in the 90–105 km altitude region is statistically significant at the 95% level. The discrepancy suggests that previous simulations of long-term thermospheric density and temperature trends<sup>3,18,19</sup> may have underestimated the enhanced cooling and contraction of the thermosphere attributable to



**Figure 4 | Average carbon profiles.** 2004–2012 average CO (red), CO<sub>2</sub> (blue) and CO<sub>x</sub> (green) VMR profiles, from ACE (solid lines) and the NCAR global mean model (dashed lines).

anthropogenic carbon emissions, because they used CO<sub>2</sub> trends from the global mean model or assumed vertically uniform relative trends. The larger observed CO<sub>2</sub> trends may therefore explain the larger observed thermospheric density decrease<sup>8</sup>. Below, we consider several possible explanations for an elevated lower thermospheric CO<sub>2</sub> trend.

The ACE-FTS trends are not a true global average, because the data represent only sunrise and sunset conditions and are weighted towards high latitudes. When the sunrise and sunset data are treated separately, the difference in trends is negligible, suggesting that there is not a strong local time dependence to the trends. There does seem to be a preference towards larger trends at lower latitudes and under high-latitude summer conditions, but those features are not statistically significant in our analysis.

A trend in atmospheric dynamics (including gravity wave activity<sup>20</sup> and stratospheric mean meridional circulation<sup>21</sup>) could potentially increase vertical mixing and advective transport, which, because of the negative gradient of the CO<sub>2</sub> (and CO<sub>x</sub>) VMR above 80 km (Fig. 4), would enhance the carbon trend in this region. There is some observational evidence that mesospheric gravity wave activity is increasing, at least in selected regions<sup>20</sup>. In the global mean model, the eddy diffusion parameter serves as a proxy for vertical mixing and advective transport; by default it has a fixed vertical profile. Imposing a trend in eddy diffusion of +15% per decade (in addition to the tropospheric CO<sub>2</sub> trend) produces a model CO<sub>x</sub> trend that is close to the observed ACE-FTS trend (Fig. 2).

A positive trend in eddy diffusion would also contribute to the anomalously low thermospheric mass densities during the 2008–2009 solar minimum<sup>22</sup>. The extra CO<sub>2</sub> drawn up from the lower mesosphere would enhance the cooling and contraction of the thermosphere, and increased downward diffusion would lower the concentration of atomic oxygen in the upper thermosphere (where it is the dominant species), thereby reducing the mass density at a fixed height. In ref. 22 it was estimated that the anomalously low densities were caused by a 14K reduction in exospheric temperature and a 12% depletion of atomic oxygen density in the lower thermosphere, compared with the previous solar minimum after accounting for the inter-minima change in solar irradiance and geomagnetic activity. In comparison, the combination of the CO<sub>2</sub> and eddy diffusion trends imposed on the global mean model produces an inter-minima exospheric temperature change of 3 K and a change of O density of 17%. Thus, these mechanisms seem to provide a plausible explanation for both the ACE-FTS CO<sub>2</sub> trends

and the anomalously low density of lower thermospheric O inferred from satellite drag measurements.

More detailed investigation of the regional effects and the contribution of atmospheric dynamics to CO<sub>2</sub> trends in the lower thermosphere will require the use of whole-atmosphere general circulation models.

In addition to tropospheric releases, another potential source of excess carbon in the upper atmosphere is *in situ* injection from the exhaust of orbital launch vehicles. For this source to contribute significantly to carbon trends, the rate of its downward diffusion must be slower than the rate of deposition, and a long-term increase in the deposition rate is probably also necessary. From the number and type of rocket launches between 2004 and 2012 (refs 23,24) we estimate a deposition of 2,700 metric tons of carbon above 90 km. In comparison, the difference between the ACE-FTS and model CO<sub>x</sub> trends between 90 and 105 km amounts to an extra 20,000 tons of carbon between 2004 and 2012. Therefore, even in the highly unlikely event that all of the carbon from space traffic exhaust has remained resident in the lower thermosphere, it cannot account for the difference between the ACE-FTS trends and the global mean model trends.

Finally, we note that the ACE-FTS CO<sub>2</sub> retrievals consider only the ground vibrational state, to minimize non-local thermodynamic equilibrium effects. At a kinetic temperature of 200 K (typical of the 90–105 km region), approximately 1.6% of the CO<sub>2</sub> population is in the excited bending state (see Supplementary Methods). A trend in the partitioning among the ground and vibrationally excited states could alias as a trend in the retrieved CO<sub>2</sub>, although it is unlikely that such an effect could explain the 10 ppm per decade difference between the ACE-FTS and global mean model trends. However, the excited states can in principle be derived from the FTS spectra, which would clarify this issue. Better knowledge of the population of each state would also illuminate the interactions among vertical mixing (which increases CO<sub>2</sub> in the mesosphere and lower thermosphere and decreases lower thermospheric O), atomic oxygen (the primary exciter of the CO<sub>2</sub> vibrational states), CO<sub>2</sub> (the primary cooling agent of the mesosphere and thermosphere) and solar ultraviolet irradiance (which photolyses CO<sub>2</sub> and O<sub>2</sub> and may drive variations in vertical mixing).

## Methods

ACE was launched in 2003 into a 74° circular orbit at an altitude of 650 km. The FTS measures high-resolution absorption spectra, covering wavelengths from 2.3 μm to 13.3 μm, during occultations of the Sun by the Earth's limb. From these spectra, height profiles of temperature, pressure and the concentrations of more than thirty species (including CO<sub>2</sub> and CO) are retrieved. Supplementary Fig. S1 shows the latitude dependence of ACE-FTS sunrise and sunset occultations as a function of day of year. Most profiles are measured at mid to high latitudes. The sampling pattern approximately repeats from year to year, owing to the 123-day orbital precession period of ACE. This sampling consistency allows us to remove the seasonal–latitude dependence of CO<sub>2</sub> and CO from the data, and thereby isolate the longer-term trends.

For this study, we used version 3.0 ACE-FTS retrievals of temperature, pressure and CO<sub>2</sub> (ref. 14). The retrievals above 65 km use small (typically 0.4 cm<sup>-1</sup> wide) regions of the CO<sub>2</sub> spectrum that have minimal interference from other molecules and contain spectral lines originating in the ground vibrational state<sup>14,17</sup>. The CO<sub>2</sub> VMR is prescribed a priori between ~25 and 60–75 km altitude; at higher altitudes, the retrieval transitions smoothly to an unconstrained, simultaneous estimation of CO<sub>2</sub> VMR and temperature<sup>14,17</sup>. The prescribed stratospheric CO<sub>2</sub> VMR includes a trend of 15 ppm per decade (ref. 17), and the influence of the imposed trends extends upward through the transition region, so the data can be used to estimate trends independently only in the lower thermosphere (above ~90 km).

For CO, we used a recently developed retrieval algorithm. The high-altitude portion of the CO retrievals employs a set of 22 lines in the CO fundamental band spanning the wavenumber range 2081–2187 cm<sup>-1</sup>, with an upper altitude of 110 km for the retrievals. Version 3.0 CO retrievals suffered problems from the assumed shape of the CO VMR profile above the highest analysed measurement: the CO VMR profile was assumed to increase rapidly with increasing altitude above 110 km, but this was inconsistent with observed CO lines in ACE-FTS spectra above 110 km. In contrast, the research product used in our analysis employed a constant CO VMR profile above the highest analysed measurement. This profile shape gives

significantly improved agreement between measured and calculated CO lines above 110 km (that is, above the altitude range employed in the CO retrieval) compared with the version 3.0 results.

We analysed ~26,500 ACE-FTS profiles from April 2004 to September 2011, using data records in which temperature and pressure were retrieved, rather than specified a priori. We excluded data from the first three months of 2004 because of spacecraft pointing problems during this early part of the mission. We first converted the profiles, which are given in geodetic altitude, to a pressure grid  $Z = \ln(P_0/P)$ , where  $P_0 = 5 \times 10^{-7}$  hPa. We then determined the seasonal–latitude dependence of the data along the sampling paths shown in Supplementary Fig. S1, by binning and averaging the data in 48 day-of-year bins (bin width 7.6 days) and treating sunrise and sunset data separately (an example is shown in Supplementary Fig. S2). Measurements more than two standard deviations away from the initial averages were omitted from the final binned average calculations. The seasonal–latitude patterns are highly consistent from year to year. We subtracted the binned averages from the VMR data in the corresponding bins to obtain a set of residuals that is nominally free of seasonal–latitude variations. We combined the sunrise and sunset residuals and averaged them in 3-month bins (excluding outliers more than two standard deviations from the initial average) to obtain a 7.5-year series of 30 average residual profiles. We estimated ACE CO<sub>x</sub> trends and their uncertainty from the time series of 3-month averages, following the method described in ref. 25, which accounts for correlated errors describable by autoregressive processes of arbitrary order (see Supplementary Methods for further information on error sources and uncertainty estimation).

The NCAR global mean model<sup>1</sup> solves the one-dimensional energy, vertical diffusion and photochemical equations governing the upper atmosphere. The lower boundary of the model, where the input CO<sub>2</sub> VMR is specified, is at  $Z = -17$  (~30 km altitude). Using version 1.2 of the model, we conducted a 1992–2012 simulation using interpolated annual global mean CO<sub>2</sub> VMRs measured in the troposphere<sup>16</sup> (<http://www.esrl.noaa.gov/gmd/ccgg/trends/global.html>) to specify lower boundary CO<sub>2</sub> variations. We added a constant offset of –13 ppm to the tropospheric time series to obtain a lower boundary CO<sub>2</sub> VMR of 372 ppm at epoch 2008.5, to approximately match the 2004–2011 ACE-FTS average at 30 km. The simulation includes solar ultraviolet irradiance and geomagnetic activity variations as parameterized by the 10.7 cm solar radio flux ( $F_{10.7}$ ) and the  $A_p$  index, respectively.

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## Author contributions

J.T.E. conceived the study, analysed the data and model output, and prepared the manuscript. M.H.S. and D.P.D. conducted the model simulations. M.H.S. calculated carbon emissions from space vehicle launches. P.F.B. and C.D.B. provided guidance on the use of the ACE-FTS retrievals. C.D.B. developed the ACE-FTS retrieval algorithms. All authors discussed the results, interpretations and implications, and contributed to the manuscript.

## Additional information

Supplementary information is available in the online version of the paper. Reprints and permissions information is available online at [www.nature.com/reprints](http://www.nature.com/reprints). Correspondence and requests for materials should be addressed to J.T.E.

## Competing financial interests

The authors declare no competing financial interests.